

Photoproduction of positive pions from hydrogen at large angles in the energy range 0.3–2.0 GeV

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Abstract. The differential cross-section for the reaction $\gamma p \rightarrow \pi^+ n$ was measured using the bremsstrahlung beam of the Bonn 2.5 GeV electron synchrotron. The pions were detected and momentum analysed in a multichannel magnet spectrometer. Data reduction resulted in 1278 measured cross-sections which are presented as energy distributions at six laboratory angles between 180° and 95° . The range of laboratory photon energies extended from 0.3 to 2.1 GeV. The statistical accuracy is better than 3 percent, the systematic error is about 5 percent. The data are compared with other experimental results and predictions of a phenomenological analysis. These cross-sections are the result of a measurement program which was started in the seventies. Further results will be reported in forthcoming publications.

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1 Introduction

Pion photoproduction up to 2 GeV incident photon energy is dominated by the formation of nucleon resonances in the s -channel. These objects have been studied extensively in πN -scattering experiments, but the electromagnetic structure of such resonant states can only be analysed through photon excitation. In order to obtain the photocouplings of the individual resonances, the photoproduction amplitudes must be evaluated by a partial-wave analysis of the experimental data. To do this without theoretical assumptions we need in principle seven independent measurements for each process, and for a proper isospin decomposition this must be done for three out of the four $\gamma N \rightarrow \pi N$ processes. The quality of such an analysis depends to a high degree on the completeness and consistency of the data used in the fit procedure.

During the last four decades a considerable amount of pion photoproduction data [1] has been accumulated, and a number of analyses [2] have been carried out. Due to the historical development of accelerator technique, the region

of the first resonance was investigated first. Later on, the measurements could be extended to the region of the second, third and fourth resonance up to about 1.4 GeV so that also in this energy range we have now abundant information. Beyond this energy the coverage in angle and energy for most of the observables is still insufficient. This applies also to the differential cross-section of positive pion photoproduction where only some measurements at small and large angles exist [1].

Measurements in the backward hemisphere are of special interest because the non-resonant background is much smaller than at forward angles, where the one-pion-exchange term gives large contributions. The excitation curve at 180° is of particular importance. At this angle the transition amplitude takes an exceptional simple form, out of four helicity amplitudes only contributions of the one having helicity 1/2 do appear. For these reasons an experiment was set up at the Bonn 2.5 GeV electron synchrotron to measure the differential cross-section for the reaction $\gamma p \rightarrow \pi^+ n$ preferably at backward angles including the laboratory angle of 180° . The aim was to take a consistent data set with high statistics, a minimum of systematic errors, and good energy resolution. The measurements presented here were made at six pion laboratory angles between 180° and 90° and photon laboratory energies covering the wide energy range from 0.3 to 2.1 GeV.

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2 Experimental method

Positive pions were photoproduced by a bremsstrahlung beam incident on a liquid hydrogen target. The produc-

tion angle Θ and the momentum p of the particles were analysed in a magnetic spectrometer. Data were taken in energy scans at six fixed laboratory angles [3]. This implied that the corresponding CM angles varied with photon energy, as shown in fig. 1. This procedure was chosen for two reasons. In the first place, a change of the angular position of the spectrometer required the break of the vacuum system between the target vessel and the detector. A continuous adjustment of the laboratory angle therefore would have led to unbearable delays in data taking. Secondly, the spectrometer collected data in 24 adjacent momentum channels simultaneously, a feature which also favoured the measurement of excitation curves at fixed laboratory angle.

Each energy scan between 0.3 GeV and 2.0 GeV required several settings of the central momentum. For each, an appropriate value for the synchrotron energy E_0 was chosen. Consecutive settings were selected so that the momentum acceptance covered half the range of the previous measurement. Thus, for each photon energy two independent measurements were carried out, however, with different parts of the momentum acceptance. This double coverage proved to be very useful for checking the consistency of the acceptance calculations and the adjustment of the system. For each combination of pion angle and momentum a minimum photon energy is needed to produce an additional pion in the reaction $\gamma p \rightarrow \pi^+ \pi^- p$ or

$\gamma p \rightarrow \pi^+ \pi^0 n$. For all our measurements, the synchrotron energy E_0 was always kept below this threshold so that multiple pion reactions could not be registered.

The background incident on the spectrometer consisted of protons, positive muons and positrons. The protons were separated from pions by time-of-flight measurements between scintillation counters. Muons originating from pion decay could not be distinguished from pions. The loss of pions due to decay and the resulting muon contamination were calculated with a Monte Carlo technique. The positron yield from pair production in hydrogen was calculated to be negligible. As an additional check, measurements with reversed polarity of the magnet were carried out.

3 Experimental arrangement

The experiment was carried out using a bremsstrahlung beam of the Bonn 2.5 GeV synchrotron [4]. Positive pions produced in a liquid-hydrogen target were identified and momentum analysed with the 0.62 GeV/ c spectrometer ATHOS. The layout of the experimental area is shown in fig. 2. The detailed aspects of the experiment are discussed below.

3.1 Photon beam and hydrogen target

The internal electron beam of the synchrotron irradiated a tungsten radiator of 1 mm diameter 10 m upstream from the hydrogen target used in this experiment. The resulting bremsstrahlung beam was collimated by the lead collimator C_1 to a rectangular cross-section with half angles of 0.9 and 0.5 mrad. The two lead scrapers C_2 and C_3 and the first sweeping magnet cleaned the beam of charged particles before entering the spectrometer area. C_2 and C_3 had slightly larger apertures than the profile of the photon beam defined by C_1 . After passing through the hydrogen target, the beam flux was measured in a Wilson-type quantameter [5]. A second cleaning magnet behind the hydrogen target removed particles scattered back from the quantameter. This was of vital importance for the measurements at 180° . The shape of the photon spectrum for each electron energy E_0 was calculated by a computer program using formulae compiled by Lublow [6]. Collimation effects were taken into account according to Lutz and Schultz [7]. The obtained spectra were estimated to be accurate to better than one percent. In order to avoid errors due to uncertainties in the shape of the spectra near the upper end, the used energy was always kept below $0.9 E_0$. The intensity profile of the photon beam was surveyed by means of silver-phosphate glasses [8] at several synchrotron energies. The alignment of the photon beam was checked by Polaroid pictures any time when the accelerator parameters were changed.

The liquid-hydrogen target was contained in a cylindrical Mylar cup 6 cm in diameter and 10 cm high. The axis

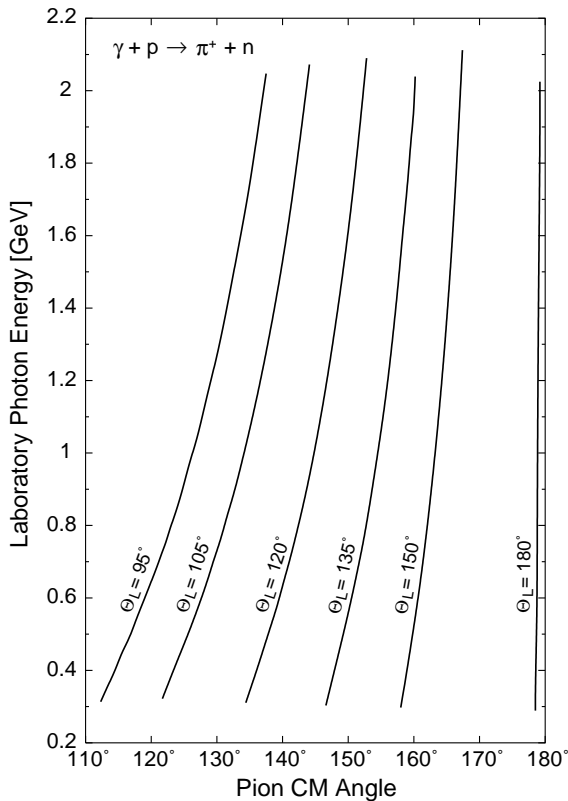


Fig. 1. Kinematics of pion photoproduction from hydrogen. The photon energies are plotted as a function of the CM angle for all laboratory angles where excitation curves were measured.

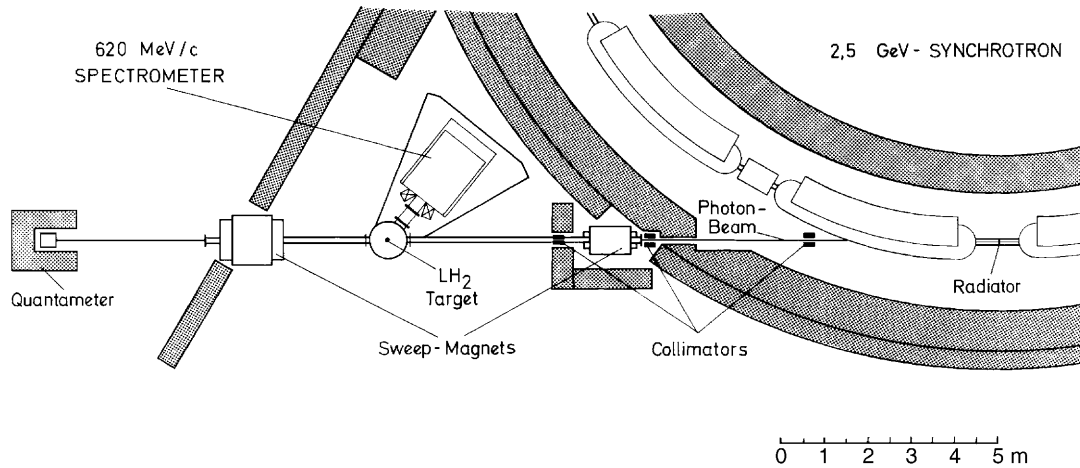


Fig. 2. Layout of the experimental area.

of the cup was vertical and perpendicular to the production plane defined by the photon beam and the spectrometer central trajectory. To prevent the residual gas from condensation on the cold walls of the target cup, an additional cell made of $7.5 \mu\text{m}$ Hostaphan foil was installed. It surrounded the hydrogen cell at a distance of 2 cm. The gap between the two foils was connected to the main vacuum system by a thin pipe only. By this provision, the formation of ice on the hydrogen cell could be kept at a negligible rate during a three weeks run. The target vessel was integrated into an extended vacuum system with a total volume of 1 m^3 . It covered the photon beam line between the two cleaning magnets and also the pion trajectories inside the spectrometer up to the front of the scintillation counter S2. The target vessel was equipped with a series of flanges to allow measurements at different angles. The vacuum chamber of the spectrometer could be connected to the target vessel at intervals of five degrees and angles between 35° and 155° . With this arrangement the background in the experimental area as well as Coulomb scattering and nuclear absorption of the pions in the spectrometer were reduced to a minimum.

3.2 The magnet spectrometer ATHOS

A side view of the $0.62 \text{ GeV}/c$ spectrometer ATHOS used in this experiment is shown in fig. 3. Charged particles emerging from the hydrogen target were deflected in the vertical direction by a wedge-shaped weak focussing magnet. The bend angle was 108° , the gap height at the center 14.5 cm. The pole pieces were 52 cm wide with a mean radius of 110 cm. For an efficient use of the field volume the pole pieces were incorporated into the vacuum chamber. The stainless-steel side walls of the chamber were skillfully soldered to the edges of the poles. The magnet was acquired from LAL Orsay where also the design and calibration had been carried out [9].

There were two connecting pipes through the iron yoke to the vacuum chamber. One of the two was oriented tangentially to the central orbit at the entrance face of the

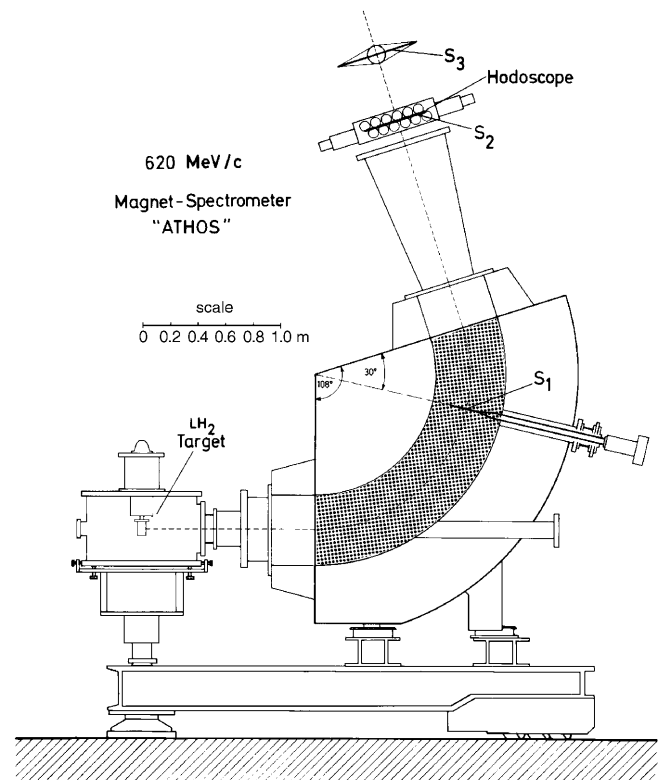


Fig. 3. Side view of the magnet spectrometer ATHOS.

magnet. Using this as entrance port for the photon beam enabled us to do the measurements at a production angle of 180° . The second one pointing radially at 30° relative to the exit face permitted an easy installation of an aperture counter inside the vacuum chamber.

At maximum induction of 1.835 T the central momentum was $0.608 \text{ GeV}/c$. The poles were shaped to give a first- and second-order field index of $n = 0.316$ and $\beta = 0.10$ which showed a small but predictable change with magnet excitation. The inclination and curvature of the effective field boundaries were corrected by shims on

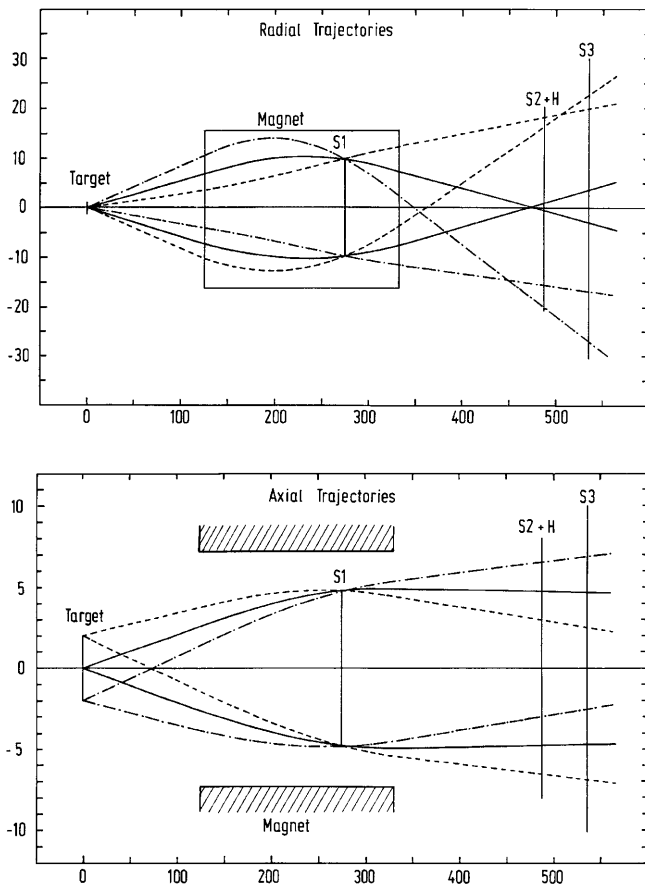


Fig. 4. Calculated trajectories through the spectrometer for initial values of horizontal and vertical angles, horizontal positions and momentum deviations.

the pole faces so that they coincided with the geometrical dimensions of the iron. The optical properties of the spectrometer are illustrated in fig. 4. In first-order optics the sector magnet was parallel-to-point focusing in the horizontal plane. In the momentum dispersing (vertical) plane there was point-to-point focusing. The hydrogen target was placed at a distance of 125 cm from the magnet entrance face. Counter S2 was located at the proper first-order focus at 144 cm behind the magnet. The position of the end counter S3 was at 192 cm. Its size was larger than necessary so that the acceptance was only determined by the dimensions of S2. The dimensions of the aperture defining counter S1 were chosen to retain all possible particle trajectories within the useful field volume and also to exclude particles scattered from the pole faces from detection. These requirements were met by a scintillator of $10 \times 20 \text{ cm}^2$ which resulted in an angle acceptance of 46 mrad in the horizontal plane, 113 mrad in the vertical plane, and an average solid angle of the spectrometer of $\Delta\Omega = 5.2 \text{ msterad}$.

Particles leaving the target with a given momentum were focused along a line in the momentum focal plane with a dispersion of 3.4 cm per percent momentum variation relative to the central trajectory. The position in this plane and thus the momentum of the particles were

measured by a momentum hodoscope. In first order the focal plane was perpendicular to the central trajectory. Second-order terms caused an inclination of 30.5° relative to the central trajectory. For technical reasons the 24 channel momentum hodoscope in front of S2 was mounted perpendicular to the central trajectory. As a consequence, the momentum resolution of the individual channels was not constant but varied between 0.5% and 1.4% FWHM. For each setting of momentum and synchrotron energy, the acceptance of all channels were determined by Monte Carlo calculations.

3.3 Electronics and on-line data processing

Pions were identified by a coincidence of the photomultiplier signals of counters S1, S2 and S3. The light produced in the scintillator of S2 was collected from two sides, and the respective PM signals S2.1 and S2.2 were mixed before they were fed into the triple coincidence $S1 \times S2.\text{mix} \times S3$ to define a pion trigger. An additional parallel logic-coincidence circuit with a 100 ns delay for signal S1 was used as a measure for the accidental-coincidence rate. A logic “OR” of the pion and the accidental-coincidence signals provided the general trigger which initialised the reading of the 24-bit hit pattern of the hodoscope as well as the time-of-flight measurement between counter S1 and S3 and also the pulse height measurement of S1, S2.1, S2.2, S2.mix and S3. For each general trigger, this information and, in addition, the contents of scalers for the quantameter rate and the time elapsed between consecutive events were transferred to an on-line computer via a CAMAC interface, and finally written on magnetic tapes for further off-line analysis. For a real-time check of the experimental parameters, the computer performed an on-line data reduction in order to display various listings, histograms and spectra. In addition, all voltages needed for the experiment were automatically checked by the computer.

4 Data collection and reduction

4.1 Experimental procedure

For each laboratory angle about 16 different settings of the spectrometer momentum and synchrotron energy were needed to cover the whole energy range. At photon energies above 0.7 GeV, we determined the cross-section for each individual hodoscope channel. Below that energy one single cross-section was calculated from the rates of two to six adjacent hodoscope channels in order to obtain a nearly constant photon-energy resolution. At energies above 1.5 GeV the energy range accepted by the spectrometer extended below the two-pion production threshold. In this case, only the data from the high-energy part of the hodoscope above this threshold was evaluated. The empty target background was measured for each setting and subtracted from the full target rates. It showed only a weak energy dependence. Up to 1 GeV photon energy

it amounted to typically 2 to 4 percent of the full target rate, and increased to 8 to 12 percent in the high-energy region. The increase of this ratio was mainly caused by the strong decrease of the differential cross-section towards higher energies. An extreme value of 46 percent was found for $\Theta = 180^\circ$ at 1.7 GeV photon energy where the pion production shows a dip.

4.2 Corrections

The net pion rates were corrected for counting losses due to nuclear absorption in the target and the counter system, for pion decay as well as for muon contamination, and counts caused by back-scattering of forward-produced pions in the hydrogen target.

The nuclear absorption correction was evaluated following the scheme of Thiessen [10]. The fraction of pions lost by nuclear interactions is shown in fig. 5. The relative uncertainty of this calculation was estimated to be about 25% which in turn caused a maximum error of the cross-section of 1.2%. At maximum momentum, 15 percent of the emitted pions decayed during their flight through the spectrometer. At the minimum detected momentum, only the fraction of $R_\pi = 0.5$ could reach counter S3. The ratio R_π as a function of the pion momentum is shown in fig. 6 together with R_μ and $R_{\pi\mu}$. The decay muons were emitted into a narrow forward cone so that some of them could continue through the counter system of the spectrometer and trigger the electronics. These events could not be distinguished from pions. The fraction R_μ of emitted pions which led to such a muon contamination was calculated in a Monte Carlo program. The total correction due to pion decay is given by $R_{\pi\mu} = R_\pi + R_\mu$. The momentum dependence of these three quantities is shown in fig. 6.

It was pointed out by d'Almagne [11] that pions originally produced in forward directions can be scattered from target nucleons into the acceptance of the spectrometer at large angles and thus give rise to a non-negligible contamination. There was again no means to separate these particles from the ones which were emitted directly into the spectrometer acceptance. The probabilities for the detection of these "phantom" pions were determined by simulating this process on the computer. Figure 7 shows the fraction of phantom pions contained in the counting rate for various laboratory angles. The uncertainty of this correction was estimated to be 20%. This contamination was especially large for the 180° measurements. The bump at 0.450 GeV is due to the large pion nucleon cross-section at this energy, whereas the peak at 1.7 GeV is again due to the dip in the photoproduction cross-section.

Effects of multiple hits in the hodoscope, gating during processing time of the on-line computer, accidentals and the influence of the finite resolutions on the cross-section were carefully investigated. The corresponding corrections turned out to be:

- multiple hits 1-3% ,
- gating 0-2.8% ,
- accidentals < 1% ,
- resolution 0-3% .

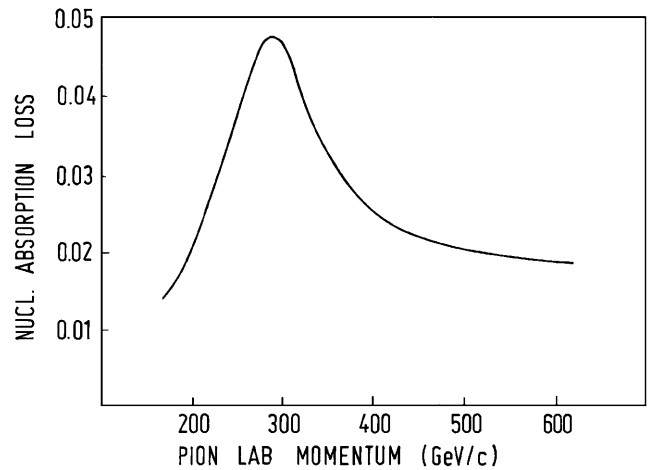


Fig. 5. Nuclear absorption correction as a function of the pion laboratory momentum.

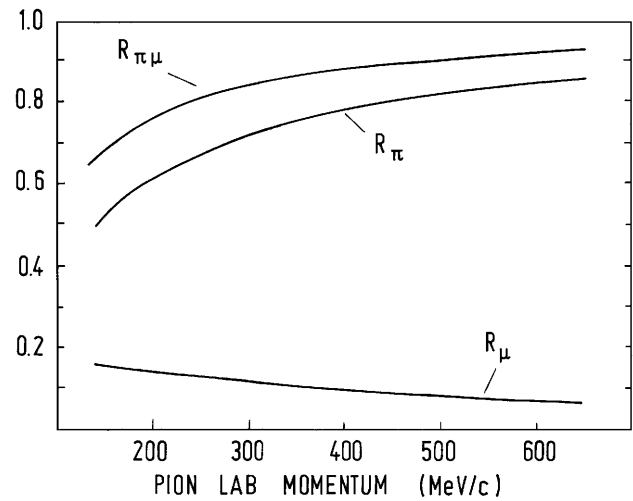


Fig. 6. Ratio of pion decay R_π , muon contamination R_μ and total rate reduction $R_{\pi\mu} = R_\pi + R_\mu$ as a function of the pion momentum.

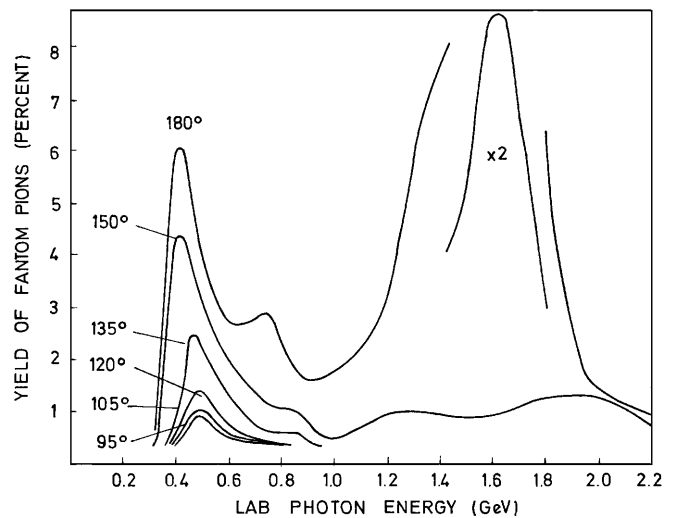


Fig. 7. Yield of back-scattered pions as a function of the laboratory photon energy.

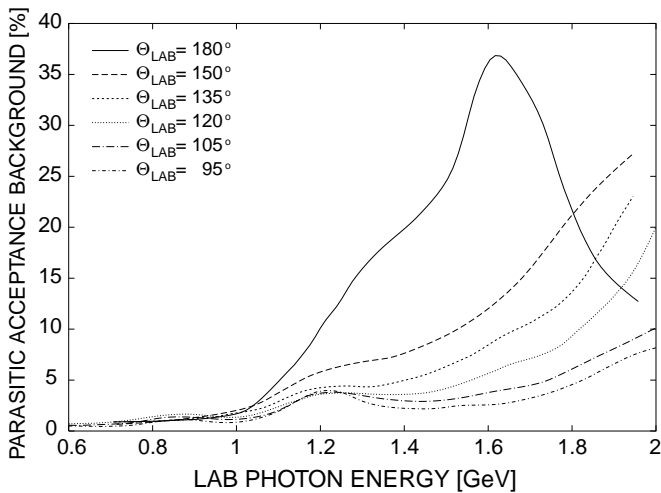


Fig. 8. Fractions of the parasitic acceptance background as a function of the laboratory photon energy.

An extensive program was carried out parallel to data taking runs to check the acceptance calculations and the momentum calibration. Various yield curves for fixed synchrotron energies E_0 were measured as a function of the spectrometer momentum. The acceptances and the calibration proved to be correct but through this procedure a new type of background was observed. When the spectrometer momentum was raised to such values that the corresponding photon energies exceeded the maximum energy E_0 , a non-vanishing yield of positive particles remained after subtraction of the empty target rates. After a large number of tests, the origin and the identity of these particles could finally be established. This intriguing background consisted out of positive pions which were produced in the hydrogen target and emitted towards the coils above the pole tips outside the regular acceptance. The pions could pass through a hidden slit between the coils, penetrate the wall of the vacuum chamber, hit the aperture counter S1, and finally reach the other counters after proper deflection by the magnetic field. These particles showed a broad momentum distribution with a center value 14% below the central spectrometer momentum. This could be determined through range measurements behind counter S4. After installation of a proper lead shielding around the coils this background disappeared.

In order to develop a suitable correction procedure for the data already collected the particle flux through this extra acceptance was investigated under various conditions by computer simulations. A fairly satisfactory description of the effect could eventually be reached by extending the momentum acceptance of each hodoscope channel towards smaller momenta by adding a linearly decaying tail to the roughly Gaussian regular acceptance function. The range and slope of this extension could be determined from the measured yield curves. Fractions of this parasitic acceptance background as a function of photon energy are shown in fig. 8. For photon energies up to 1.0 GeV, the calculated corrections were less than two percent of the measured pion rates at all angles. At higher energies we

observed a strong dependence on the pion angle. For the excitation curves at 95° and 105° the correction amounted to 5% at 1.8 GeV and 8%–10% at 2.0 GeV. For the measurements at laboratory angles of 120°, 135° and 150° we obtained corrections of 4%–8% at $E_\gamma = 1.5$ GeV with a steep increase to 20%–28% at 2 GeV. At $\Theta = 180^\circ$ this correction was exceptionally large. It rose almost linearly to 35% at 1.6 GeV, where the cross-sections itself shows its minimum, and then dropped down again to 10% at $E_\gamma = 2$ GeV. Because of the somewhat simplified assumptions and some uncertainties in the determination of the parameters, the relative uncertainty of these corrections was estimated to be 20 percent.

5 Results and discussion

5.1 Results of this experiment

The cross-sections measured in this experiment are plotted in figs. 9 to 14, and listed in table 1 as excitation curves at fixed laboratory angles. The errors quoted with the cross-sections include counting statistics and the uncertainties of the parasitic acceptance correction but no other systematic errors. The photon-energy resolution $\Delta k/k$ depends on the pion laboratory angle and the photon energy. It varied from 1.5% (FWHM) at $\Theta = 180^\circ$ in the first resonance region up to 8.5% (FWHM) at 2.0 GeV for $\Theta = 95^\circ$. The angular resolution was typically 1.7° (FWHM) and reached a maximum of 2.3° (FWHM) at $\Theta = 180^\circ$. All excitation curves show a general decrease with energy superimposed by maxima around 0.3 GeV, 0.7 GeV, 1.0 GeV and 1.4 GeV known as the first, second, third, and fourth resonance. In order to display adequately the structures in the cross-sections at all energies, the data around the first resonance were reduced to 1/4. Above the third resonance the cross-sections are in addition plotted after multiplication by 4.

Measurements at backward angles are especially useful for the study of the resonances in the low partial waves which are, in general, veiled by the dominant resonance contributions in the higher waves. By angular momentum conservation at an angle of 180°, all helicity amplitudes vanish but H_4 which corresponds to an initial helicity 1/2. As was first observed by Walker [12], the second and third resonances occur almost entirely in the amplitudes corresponding to initial helicity 3/2, and thus show little effect on the backward cross-section. In addition, the t -channel term of the electric Born approximation which contains high partial waves also vanishes at 180°. The remaining electric and magnetic Born terms only affect the lowest partial waves.

Looking at the excitation curve at $\Theta = 180^\circ$ (fig. 9) we notice the following features. Above 0.3 GeV the data show the steep slope of the high-energy tail of the first resonance up to about 0.5 GeV. In the following range, up to 1 GeV, the cross-section in general decreases slowly and then drops rapidly towards higher energies. After a shallow minimum around 1.65 GeV the cross-section

Pion-Lab.-Angle: 95 Degree						Pion-Lab.-Angle: 95 Degree						Pion-Lab.-Angle: 95 Degree					
k_{lab} [GeV]	W_{CM} [GeV]	Θ_{CM} [deg.]	$d\sigma/d\Omega$ [$\mu\text{b}/\text{sr}$]	Error [$\mu\text{b}/\text{sr}$]	#	k_{lab} [GeV]	W_{CM} [GeV]	Θ_{CM} [deg.]	$d\sigma/d\Omega$ [$\mu\text{b}/\text{sr}$]	Error [$\mu\text{b}/\text{sr}$]	#	k_{lab} [GeV]	W_{CM} [GeV]	Θ_{CM} [deg.]	$d\sigma/d\Omega$ [$\mu\text{b}/\text{sr}$]	Error [$\mu\text{b}/\text{sr}$]	#
0.758	1.517	122.3	5.141	0.128	9	1.056	1.692	127.2	2.669	0.054	11	1.782	2.055	135.4	0.230	0.007	15
0.758	1.517	122.3	4.987	0.139	8	1.067	1.698	127.3	2.509	0.051	11	1.807	2.067	135.5	0.206	0.007	15
0.765	1.522	122.4	4.858	0.124	9	1.077	1.703	127.4	2.204	0.049	11	1.836	2.080	135.8	0.201	0.007	15
0.766	1.522	122.4	4.828	0.132	8	1.090	1.711	127.7	1.887	0.045	11	1.838	2.081	135.8	0.190	0.007	16
0.771	1.526	122.5	4.737	0.123	9	1.101	1.717	127.8	1.700	0.043	11	1.864	2.092	136.0	0.176	0.006	15
0.772	1.526	122.5	4.538	0.127	8	1.113	1.723	128.0	1.481	0.040	11	1.865	2.093	136.0	0.171	0.007	16
0.778	1.530	122.6	4.559	0.119	9	1.126	1.730	128.2	1.257	0.037	11	1.891	2.105	136.2	0.153	0.007	16
0.780	1.531	122.7	4.329	0.125	8	1.138	1.737	128.3	1.067	0.035	11	1.922	2.118	136.5	0.127	0.006	16
0.784	1.533	122.8	4.211	0.114	9	1.138	1.737	128.3	1.149	0.036	12	1.953	2.132	136.8	0.115	0.006	16
0.787	1.535	122.8	4.138	0.125	8	1.149	1.743	128.4	0.953	0.034	12	1.986	2.146	137.0	0.109	0.006	16
0.791	1.538	122.9	4.046	0.116	9	1.151	1.744	128.5	0.927	0.034	11	2.018	2.160	137.3	0.100	0.005	16
0.798	1.542	123.0	4.045	0.112	9	1.162	1.750	128.6	0.931	0.032	12	2.047	2.173	137.5	0.098	0.005	16
0.805	1.546	123.1	3.822	0.108	9	1.173	1.755	128.8	0.855	0.030	12						
0.811	1.550	123.2	3.831	0.104	10	1.186	1.762	129.0	0.721	0.030	12						
0.812	1.551	123.2	3.537	0.108	9	1.199	1.769	129.0	0.693	0.028	12						
0.817	1.554	123.3	3.685	0.101	10	1.212	1.776	129.3	0.659	0.028	12						
0.821	1.556	123.4	3.489	0.102	9	1.226	1.784	129.5	0.713	0.028	12						
0.824	1.558	123.5	3.530	0.102	10	1.240	1.791	129.6	0.652	0.027	12						
0.828	1.560	123.5	3.296	0.101	9	1.255	1.799	129.9	0.583	0.026	12						
0.830	1.561	123.6	3.493	0.098	10	1.269	1.806	130.0	0.628	0.026	12						
0.836	1.565	123.6	2.931	0.098	9	1.285	1.814	130.2	0.647	0.026	12						
0.837	1.566	123.7	3.249	0.095	10	1.301	1.823	130.4	0.683	0.024	12						
0.844	1.570	123.8	3.194	0.098	10	1.317	1.831	130.6	0.682	0.025	12						
0.844	1.570	123.8	3.025	0.093	9	1.318	1.831	130.6	0.689	0.024	13						
0.851	1.574	123.9	3.071	0.094	10	1.332	1.838	130.7	0.709	0.025	13						
0.859	1.579	124.1	3.123	0.089	10	1.333	1.839	130.8	0.684	0.024	12						
0.866	1.583	124.2	3.078	0.088	10	1.349	1.847	131.0	0.711	0.024	13						
0.873	1.587	124.3	3.201	0.093	10	1.364	1.855	131.1	0.694	0.024	13						
0.881	1.592	124.4	3.063	0.091	10	1.381	1.863	131.3	0.705	0.024	13						
0.889	1.596	124.6	3.061	0.089	10	1.399	1.872	131.5	0.686	0.023	13						
0.897	1.601	124.7	3.088	0.088	10	1.415	1.880	131.7	0.702	0.023	13						
0.905	1.606	124.8	3.057	0.091	10	1.435	1.890	132.0	0.668	0.023	13						
0.914	1.611	125.0	3.159	0.089	10	1.452	1.899	132.1	0.656	0.022	13						
0.923	1.616	125.1	2.961	0.089	10	1.472	1.909	132.3	0.661	0.023	13						
0.932	1.622	125.3	3.133	0.089	10	1.494	1.919	132.6	0.609	0.021	13						
0.941	1.627	125.4	3.358	0.090	10	1.506	1.925	132.7	0.649	0.018	14						
0.950	1.632	125.5	3.304	0.087	10	1.513	1.929	132.8	0.614	0.020	13						
0.959	1.637	125.7	3.361	0.092	10	1.521	1.933	132.8	0.592	0.017	14						
0.969	1.643	125.8	3.549	0.063	11	1.541	1.942	133.0	0.575	0.017	14						
0.969	1.643	125.8	3.389	0.093	10	1.561	1.952	133.2	0.538	0.016	14						
0.978	1.648	126.0	3.467	0.062	11	1.582	1.962	133.5	0.509	0.015	14						
0.979	1.648	126.0	3.578	0.093	10	1.606	1.973	133.7	0.500	0.015	14						
0.986	1.652	126.1	3.576	0.062	11	1.627	1.983	133.9	0.436	0.015	14						
0.990	1.655	126.2	3.415	0.093	10	1.652	1.995	134.1	0.418	0.014	14						
0.996	1.658	126.3	3.565	0.063	11	1.675	2.006	134.4	0.364	0.014	14						
1.005	1.663	126.4	3.579	0.062	11	1.682	2.009	134.4	0.386	0.009	15						
1.015	1.669	126.6	3.499	0.062	11	1.700	2.018	134.6	0.347	0.013	14						
1.025	1.674	126.8	3.393	0.060	11	1.722	2.030	134.8	0.345	0.008	15						
1.035	1.680	129.8	3.059	0.059	11	1.724	2.030	134.8	0.319	0.008	15						
1.044	1.685	127.0	2.874	0.055	11	1.754	2.043	135.1	0.270	0.008	15						

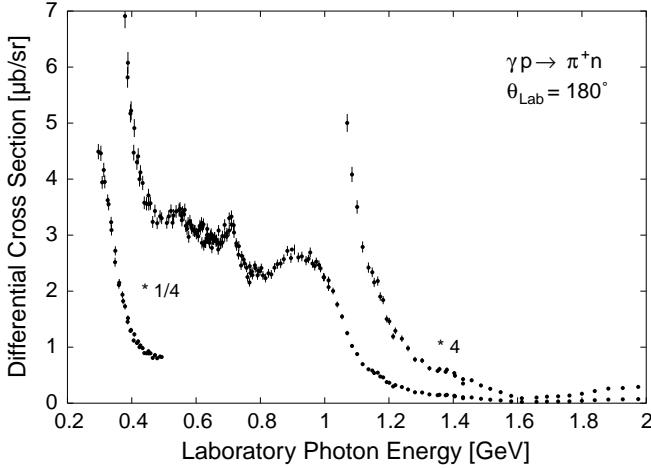


Fig. 9. Differential cross-section for $\gamma p \rightarrow \pi^+ n$ at a laboratory pion angle of $\Theta = 180^\circ$ as a function of the laboratory photon energy. The data around the first resonance are reduced to $1/4$ and above the third resonance the cross-sections are in addition plotted after multiplication by 4.

increases again. Centered at the threshold for the η -photoproduction $\gamma p \rightarrow \eta p$ at 0.71 GeV there appears a sharp spike. The following dip is complementary to the η -photoproduction cross-section which has its maximum around 0.8 GeV and then drops rapidly to 20% of the maximum value above 0.9 GeV. In a simple picture, the observed phenomenon can be explained by a strong ex-

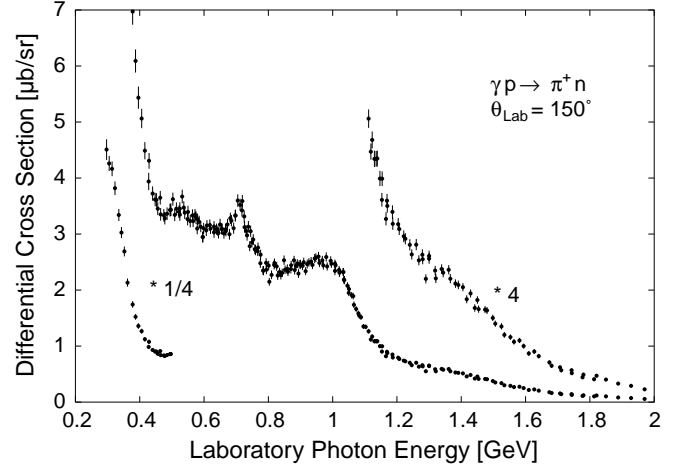


Fig. 10. Differential cross-section for $\gamma p \rightarrow \pi^+ n$ at a laboratory pion angle of $\Theta = 150^\circ$ as a function of the laboratory photon energy. The data around the first resonance are reduced to $1/4$ and above the third resonance the cross-sections are in addition plotted after multiplication by 4.

citation of the $S_{11}(1535)$ resonance which above the η -production threshold at 0.71 GeV predominantly decays into $p + \eta$, and thus, by unitarity, causes a discontinuity in the pion-nucleon channel. This η -cusp in $\gamma p \rightarrow \pi^+ n$ was reported earlier by Hand and Schaerf [13]. A more detailed investigation of this effect was given in ref. [14].

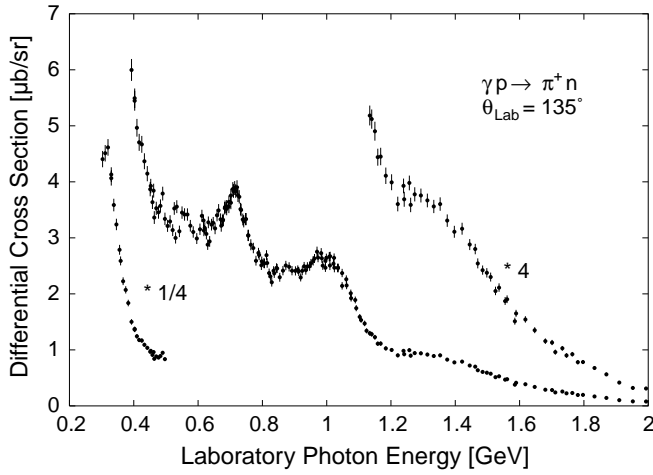


Fig. 11. Differential cross-section for $\gamma p \rightarrow \pi^+ n$ at a laboratory pion angle of $\Theta = 135^\circ$ as a function of the laboratory photon energy. The data around the first resonance are reduced to $1/4$ and above the third resonance the cross-sections are in addition plotted after multiplication by 4.

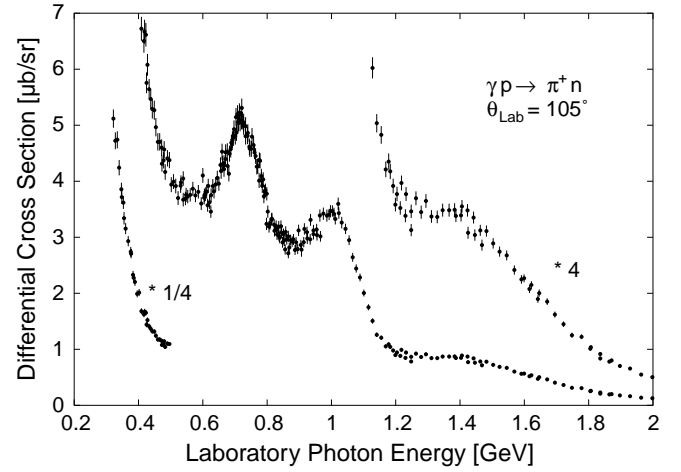


Fig. 13. Differential cross-section for $\gamma p \rightarrow \pi^+ n$ at a laboratory pion angle of $\Theta = 105^\circ$ as a function of the laboratory photon energy. The data around the first resonance are reduced to $1/4$ and above the third resonance the cross-sections are in addition plotted after multiplication by 4.

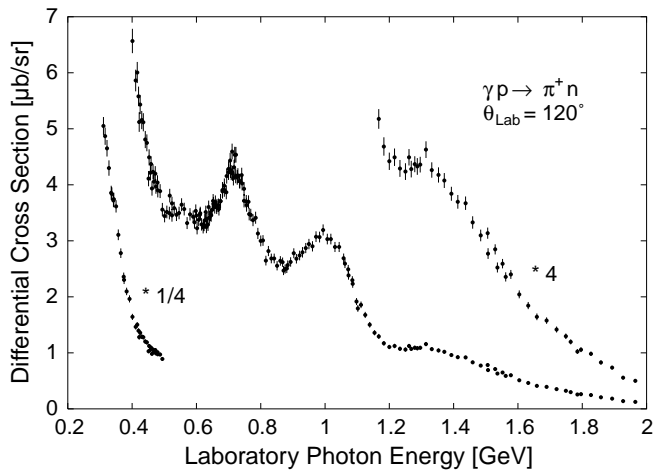


Fig. 12. Differential cross-section for $\gamma p \rightarrow \pi^+ n$ at a laboratory pion angle of $\Theta = 120^\circ$ as a function of the laboratory photon energy. The data around the first resonance are reduced to $1/4$ and above the third resonance the cross-sections are in addition plotted after multiplication by 4.

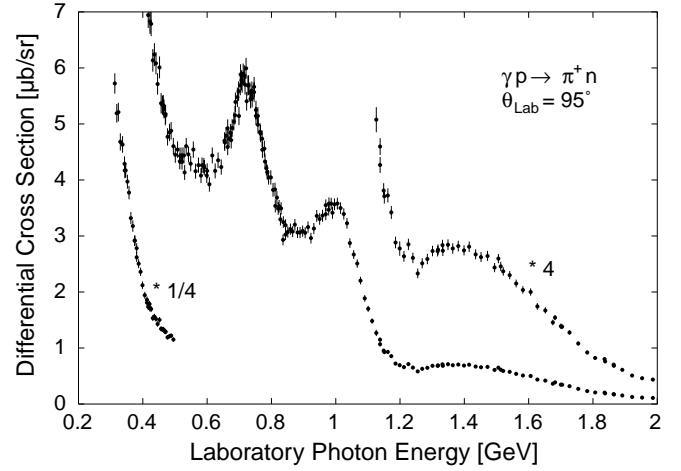


Fig. 14. Differential cross-section for $\gamma p \rightarrow \pi^+ n$ at a laboratory pion angle of $\Theta = 95^\circ$ as a function of the laboratory photon energy. The data around the first resonance are reduced to $1/4$ and above the third resonance the cross-sections are in addition plotted after multiplication by 4.

Changing to a laboratory angle of 150° (fig. 10) we find a similar behaviour of the measured cross-section up to about 1.1 GeV, except for the η -cusp region where the observed structure is shifted to higher cross-section values because of the build-up of the second resonance. The most striking feature is the steep increase with decreasing angle at energies around 1.25 GeV. At the laboratory angles of 135° , 120° , 105° and 95° (figs. 11 to 14) we find a slow increase of the cross-section at the first resonance. The second resonance is growing much stronger towards smaller angles. At all angles the opening of the η -channel shows up as an almost linear decrease of the cross-section above 0.71 GeV on top of the second resonance. We notice a moderate growth of the cross-section at the position of the third resonance at 1.0 GeV. At photon energies

of 1.35 GeV the cross-section increases strongly towards smaller angles, reaches a maximum at a laboratory angle of 120° , and then drops again towards 95° .

5.2 Comparison with existing data and the recent SAID analysis

At photon energies below 2 GeV a large number of experiments has been carried out over the last four decades. In the following, we will restrict our discussion to the measurements which were accepted to the data base for the SAID analysis carried out by Arndt, Strakovsky and Workman [15]. Our results in the region of the first resonance at six different laboratory angles are plotted in

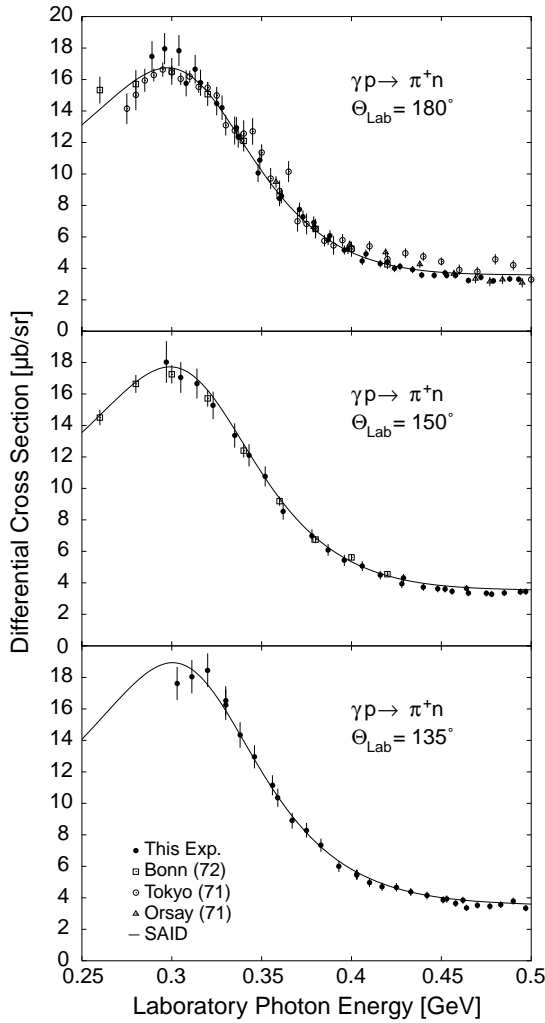


Fig. 15. The results of this experiment in the region of the first resonance in comparison with the SAID analysis and data from Bonn [16], Tokyo [18] and Orsay [17].

figs. 15 and 16 together with the curves of the recent SAID analysis. (The error bars given here also include the systematic uncertainties.) In addition, the results of measurements from Tokyo, Orsay, Mainz and Bonn are shown. Within the experimental uncertainties our data agree quite well with the SAID results. This is also true for the earlier Bonn results [16]. At the laboratory angle of 180° there are data from two dedicated experiments carried out at Orsay [17] and Tokyo [18]. Around the maximum of the first resonance the Tokyo values in general agree with the SAID curves. Above 400 MeV, however, we notice larger deviations. With a few exceptions, the Orsay data are in agreement with the SAID analysis.

At smaller laboratory angles there are further contributions from Tokyo, Mainz and Bonn. The Tokyo results [19] for $\Theta_{\text{Lab}} = 120^\circ$ follow within the quoted uncertainties the SAID curve with a few exceptions. At $\Theta_{\text{Lab}} = 120^\circ$ and the two smaller laboratory angles there are data from two Mainz experiments [20,21]. Along the high-energy tails of the first resonance these values follow

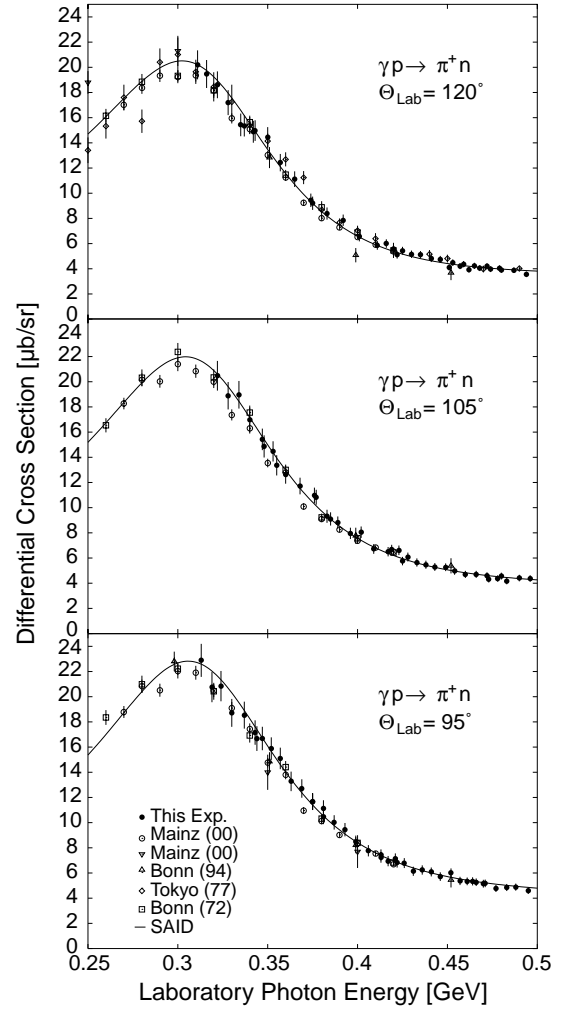


Fig. 16. The results of this experiment in the region of the first resonance in comparison with the SAID analysis and data from Mainz [20], Mainz [21], Bonn [16], Tokyo [19] and Bonn [22].

closely the SAID curves. Around the resonance maximum there seems to be a slight difference in normalisation. The Bonn data [22] at 400 MeV and 450 MeV agree very well for $\Theta_{\text{Lab}} = 105^\circ$ and $\Theta_{\text{Lab}} = 95^\circ$, whereas at $\Theta_{\text{Lab}} = 120^\circ$ there are sizeable deviations. A comparison of our data at photon energies between 450 and 1250 MeV with the SAID analysis is given in figs. 17 to 22. The error bars shown here also include the systematic uncertainties. The data from Orsay, Tokyo, Cornell and MIT are also plotted.

In general, our data follow the course of the SAID excitation curves very well, though there are also a few local deviations. The measurements at the laboratory angle of 180° are shown in fig. 17. At the low-energy side, the Tokyo data [18] show considerable deviations, but good agreement above 700 MeV. The Orsay [17] results indicate a similar behaviour. At photon energies above 660 MeV there are data from a further dedicated 180° measurement carried out at Cornell [23]. In the beginning they agree with the SAID results, but above 800 MeV are large deviations. At the smaller laboratory angles of $\Theta_{\text{Lab}} = 135^\circ$ and $\Theta_{\text{Lab}} = 95^\circ$ the Tokyo data [19] are

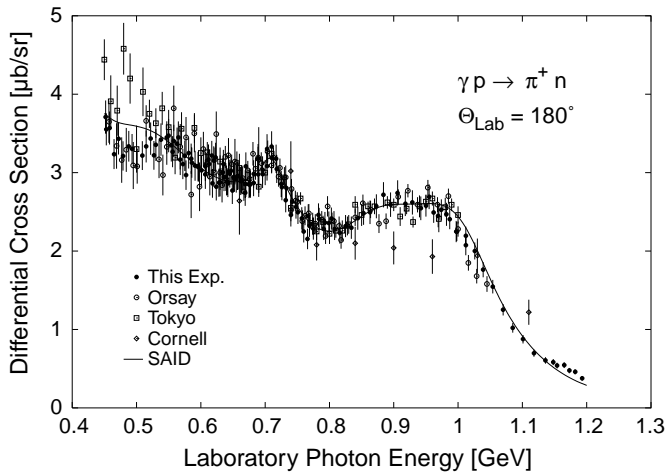


Fig. 17. The results of this experiment in the medium energy range at the laboratory angle of 180° compared with the SAID analysis and data from Orsay [17], Tokyo [18] and Cornell [23].

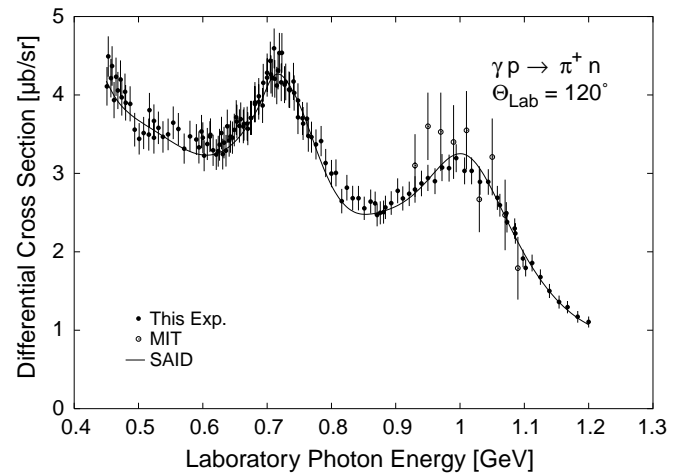


Fig. 20. The results of this experiment in the medium energy range at a laboratory angle of 120° compared with the SAID analysis and data from MIT [19].

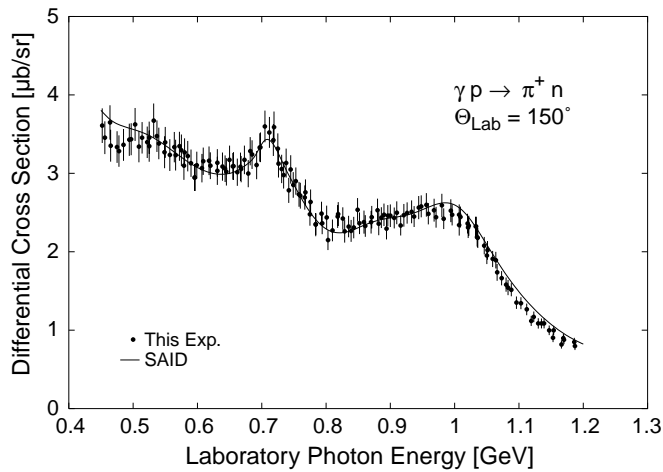


Fig. 18. The results of this experiment in the medium energy range at the laboratory angle of 150° compared with the SAID analysis.

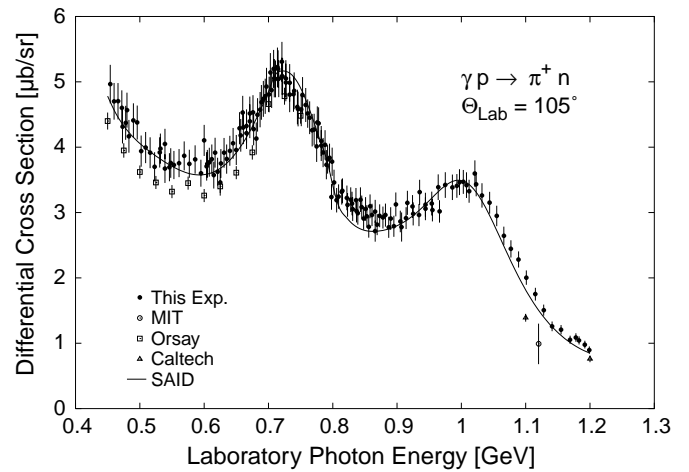


Fig. 21. The results of this experiment in the medium energy range at a laboratory angle of 105° compared with the SAID analysis and data from Orsay [24].

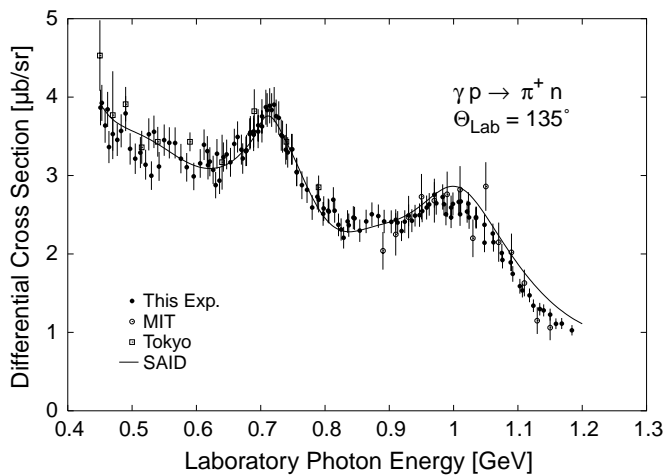


Fig. 19. The results of this experiment in the medium energy range at the laboratory angle of 135° compared with the SAID analysis and data from MIT [25] and Tokyo [19].

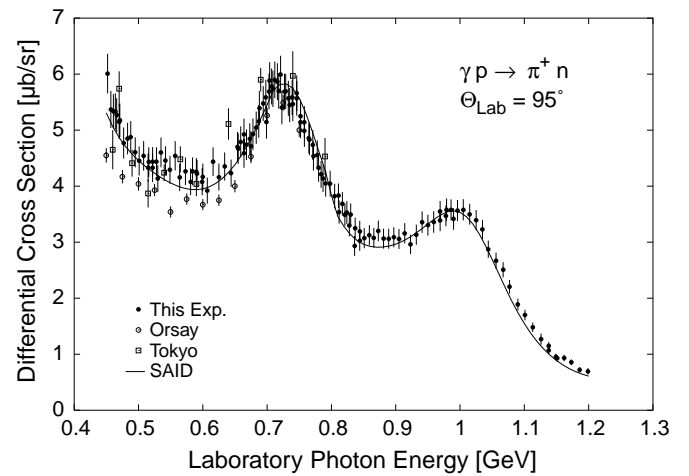


Fig. 22. The results of this experiment in the medium energy range at a laboratory angle of 95° compared with the SAID analysis and data from Orsay [24] and Tokyo [19].

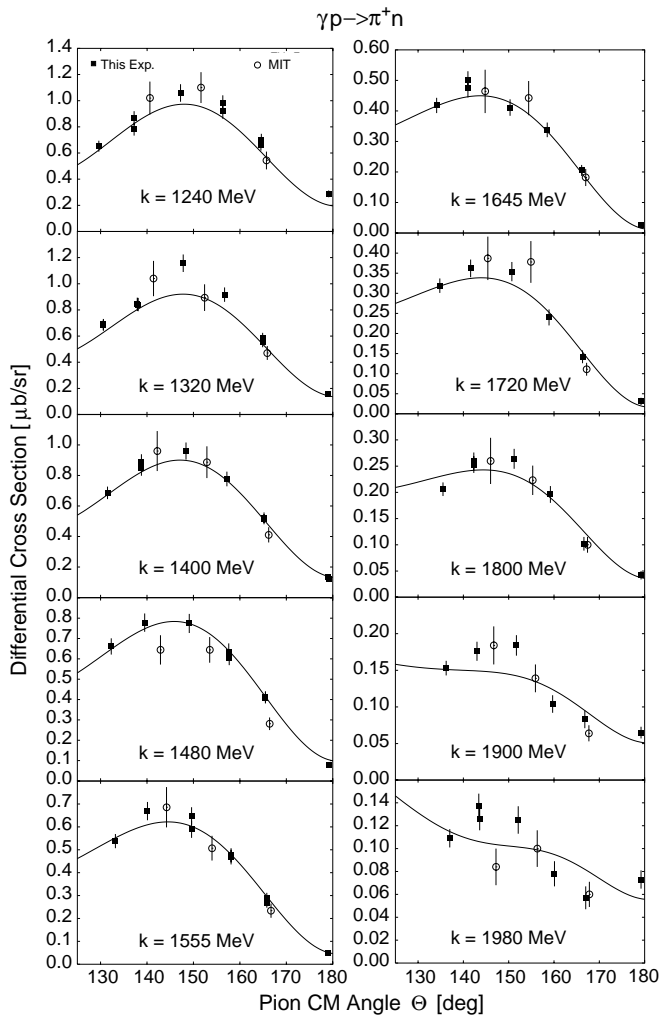


Fig. 23. CM angular distributions at backward angles for photon energies between 1.240 MeV (GeV) and 1.980 MeV (GeV). The results of this experiment are plotted together with the data from MIT [25]. The curves show the results of the SAID analysis.

systematically above the SAID curves, whereas the Orsay values [24] at $\Theta_{\text{Lab}} = 105^\circ$ and $\Theta_{\text{Lab}} = 95^\circ$ are always below the analysis. The values of MIT [25] at $\Theta_{\text{Lab}} = 135^\circ$ and $\Theta_{\text{Lab}} = 120^\circ$ are in agreement with the SAID curves, but they show partly large uncertainties. The MIT group measured three excitation curves at three fixed laboratory angles in the backward hemisphere above photon energies of 0.900 GeV. Part of their data between 1.2 GeV and 1.98 GeV are plotted as angular distributions at ten fixed photon energies in fig. 23 together with our cross-sections at the corresponding energies. Again the curves give the results of the SAID analysis. Within the experimental uncertainties the two data sets agree with the analysis, except at the two highest energies where the combined

data sets indicate a different shape of the angular distributions.

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